

# Jet Physics At HERA

Roman Kogler

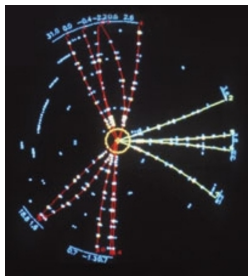
Max-Planck-Institute for Physics  
Munich

Hamburg Student Seminar

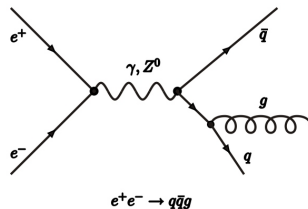
June 26<sup>th</sup>, 2008



# Why Jets?



3-jet event at TASSO,  
discovery of the gluon in 1979



a gluon-bremsstrahlung diagram  
in  $e^+e^-$  annihilation

**Detecor:** Jets are regions of high energy density

**Hadrons:** Jets are collimated sprays of hadrons

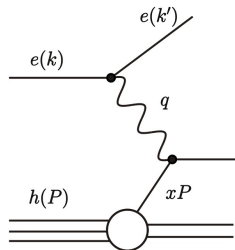
**Partons:** Jets are formations of close-by partons (jet  $\neq$  parton)

*Find a mapping between these levels!*

# A Little Bit Of QCD



# DIS Kinematics



$q^\mu = k^\mu - k'^\mu$ , momentum transfer

$Q^2 = -q^2$ , virtuality

$x = \frac{Q^2}{2P \cdot q}$ , Bjorken scaling variable

$y = \frac{P \cdot q}{P \cdot k}$ , inelasticity

At HERA,  $E_p = 920$  (820) GeV and  $E_e = 27.6$  GeV  
center-of-mass energy  $\sqrt{s} \approx 320$  (300) GeV

Beyond the DIS Born diagram having some hard scale in the event,  
e.g.  $Q^2 \gtrsim 3 \text{ GeV}^2$ , **pQCD is applicable**

Testing ground and precision measurements of QCD

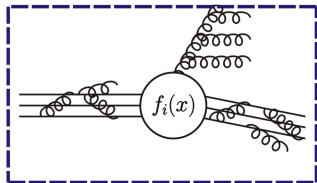
# Factorization Theorem In DIS

$$d\sigma_{ep} = \sum_{i=q\bar{q}g} \int_0^1 dx f_i(x, \mu_F^2; \alpha_s) d\hat{\sigma}_i(xP, \alpha_s, \mu_R, \mu_F)$$



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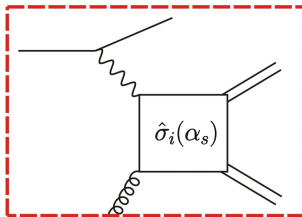
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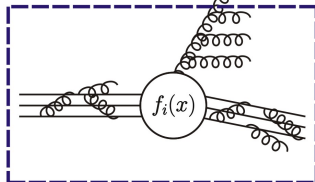
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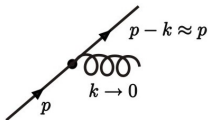
$d\hat{\sigma}_i(xP, \alpha_s, \mu_R, \mu_F)$  : partonic hard cross section, calculable in pQCD, IR safe quantity, independent of long-range effects, i.e. soft ( $p_T < \mu_F$ ) physics



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# Calculation Of The Short Distance (Hard) Part

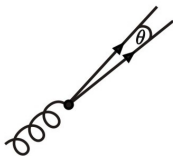
Already at one-loop  $\mathcal{O}(\alpha_s)$  the hard **partonic** cross section is divergent:



divergent for  $k \rightarrow 0$ :

$$\int_0 \frac{dk}{k} \rightarrow \infty$$

IR divergency



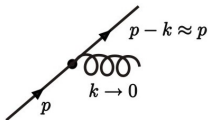
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collinear divergency

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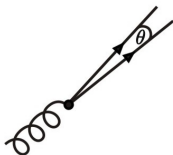
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**Solution:** define IR safe quantities which do not change under soft emissions or collinear splittings. Also ensures that the total cross section is finite.

# Jet Algorithms



# Requirements On A Jet Algorithm

From the theoretical considerations, a jet algorithm should exhibit

- infrared safety
- collinear safety
- factorizability - separation from the beam remnant(s)
- small renormalization scale dependence

These are required for meaningful (finite) fixed-order calculations (with small errors).



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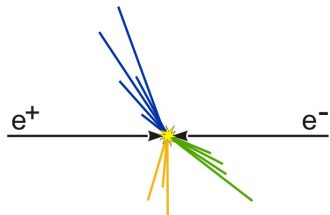
Some experimental/phenomenological aspects:

- approximate insensitivity to the soft underlying event, pileup
- fast running time
- close correspondance between final state partons and hadronic final state (small hadronization corrections)

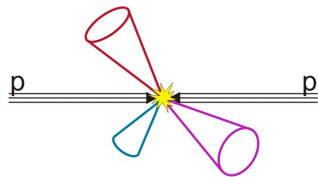
These are required for the application of a jet algorithm in an analysis.



# Historical Development

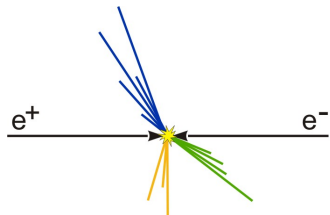


Work with angles, angular 'distance' lead to successive recombination (clustering) algorithms



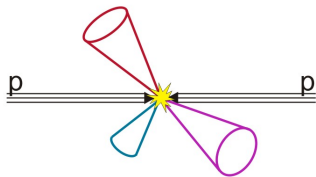
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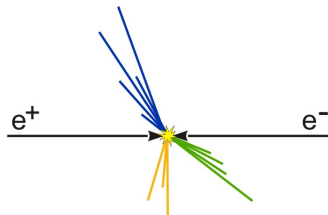
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They work because QCD modifies the energy flow on small scales

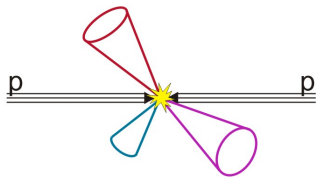
# Historical Development



Work with angles, angular 'distance' lead to successive recombination (clustering) algorithms

They work because the proximity measure is closely related to QCD divergencies

simplicity and modest hadronization corrections



Work with coarse regions of high energy density, lead to cone-type algorithms

They work because QCD modifies the energy flow on small scales

geometrical regularity and lower sensitivity to non-perturbative underlying event and pile-up

# A Modern Clustering Algorithm

## The longitudinally invariant $k_T$ algorithm

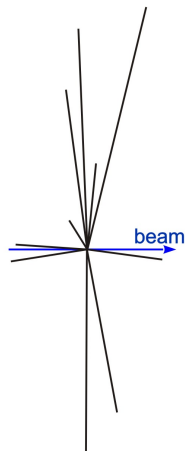
distance between objects  $i$  and  $j$ :

$$d_{ij} = \min(k_{T,i}, k_{T,j}) \cdot \sqrt{(\Delta\eta_{ij})^2 + (\Delta\phi_{ij})^2/R_0}$$

beam distance:

$$d_{iB} = k_{T,i}$$

make a list of all possible  $d_{ij}$  and  $d_{iB}$   
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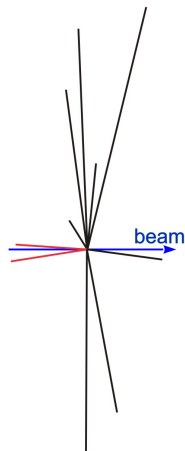
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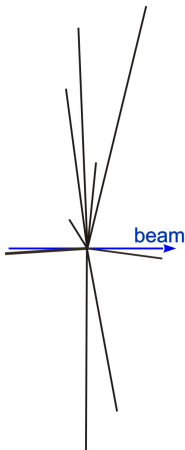
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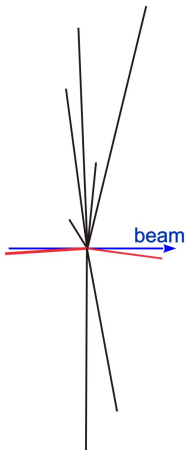
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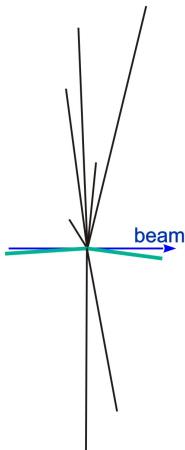
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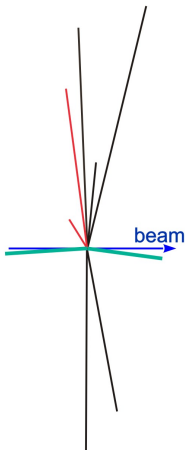
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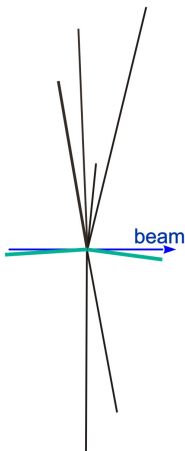
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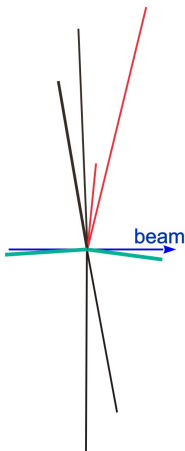
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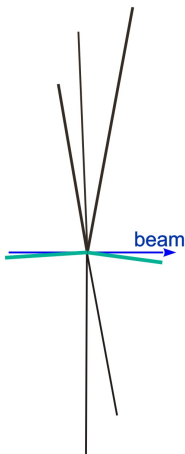
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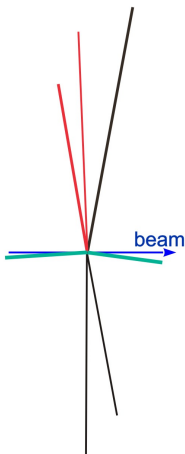
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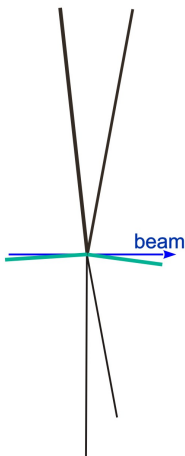
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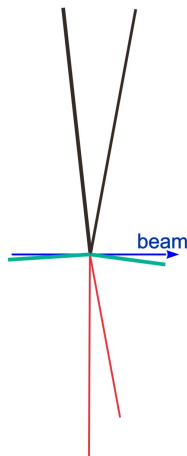
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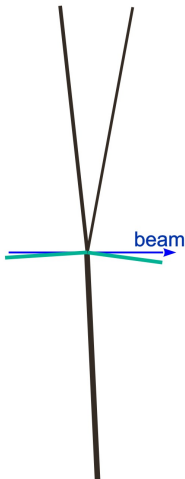
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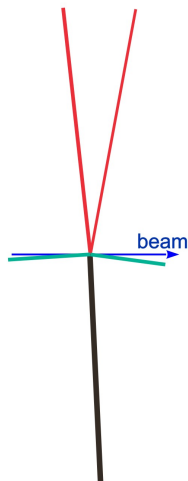
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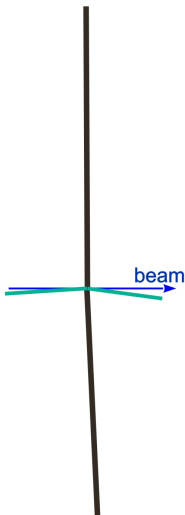
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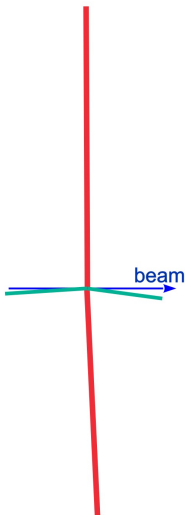
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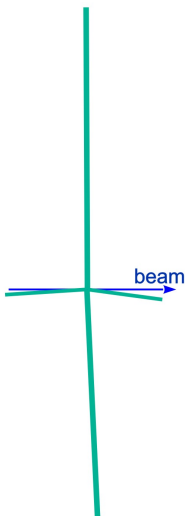
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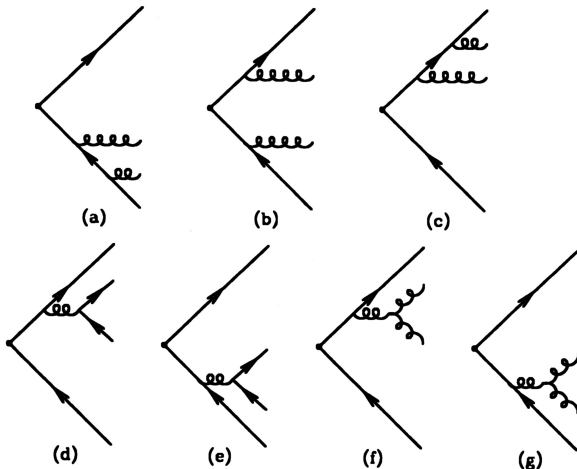
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Two classes of tests:

- 1) Test jet-production cross sections in various regions of phase space.
- 2) More intimate testing: jet shapes, radius dependence, heavy flavour jets, ...



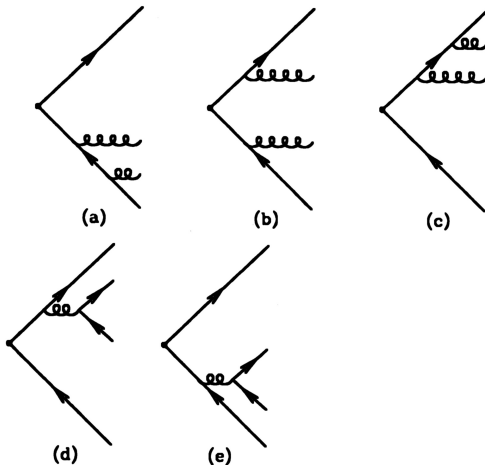
# Is QCD Non-Abelian?



All LO diagrams for  $e^+e^- \rightarrow q\bar{q}q\bar{q}$ ,  $q\bar{q}gg$  contribution to 4-jet production in  $e^+e^-$  annihilation

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Abelian QCD', [U(1)]<sup>3</sup> has less  $q\bar{q}gg$  contributions to the 4-jet cross section

# Is QCD Non-Abelian?

A. Ali et al., Phys. Lett. B82 (1979) 285:

$$d\sigma^{4jet} = \left(\frac{\alpha_s}{4\pi}\right)^2 [C_F T_F n_f \mathcal{A} + C_F^2 \mathcal{B} + C_F C_A \mathcal{C}]$$



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## Strategy:

Fix  $\alpha_s$  by 3-jet rate and calculate  $d\sigma^{4jet}$  for both theories

Experimental test is difficult because the absolute normalization changes only slightly

Find different observables!



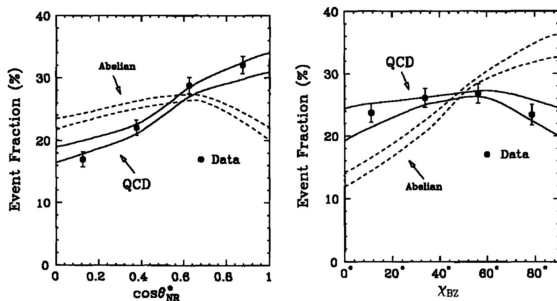
# Is QCD Non-Abelian?

L3 collaboration (LEP): Phys. Lett. B 248 (1990), 227:  
Measure the relative orientation of the planes containing the primary  $q\bar{q}$  jets and the secondary  $q\bar{q}$  or  $gg$  jets



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SU(3) is clearly favored compared to  $[U(1)]^3$

Nachtmann-Reiter angle  $\theta_{NR}^*$  is the angle between the momentum vector differences of the primary and secondary jets

Bengtsson-Zerwas angle  $\chi_{BZ}$  is the angle between the planes of the primary and secondary jets



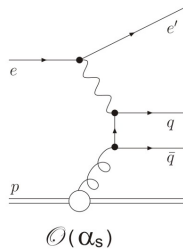
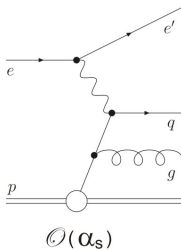
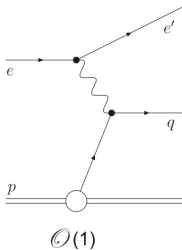
# Inclusive Jets At HERA

$$d\sigma_{ep}^{jet} = \sum_{i=q\bar{q}g} \int_0^1 dx f_i(x, \mu_F^2; \alpha_s) d\hat{\sigma}_i^{jet}(xP, \alpha_s, \mu_R, \mu_F) \cdot C_{had}$$



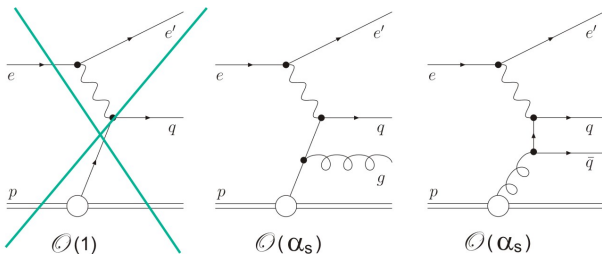
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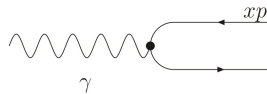


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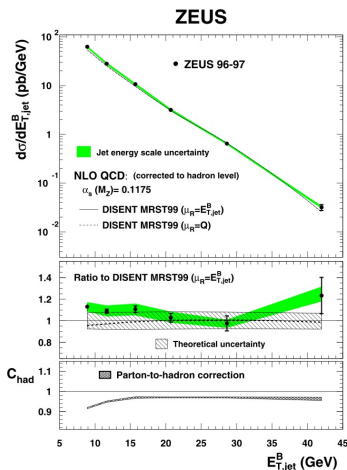


To exclude QPM  $\mathcal{O}(\alpha_s^0)$  diagrams,  
boost to the Breit frame and cut on  $P_T^{jet}$



# Inclusive Jets At HERA

ZEUS, Physics Letters B 547 (2002) 164-180:



PDF was determined from global fits without any HERA jet-data - independent test of its validity

calculation was corrected to hadron level ( $C_{had}$ ) and for  $Z^0$  exchange

fixed-order calculation  $\mathcal{O}(\alpha_s^2)$  describes the data very well over more than 3 orders of magnitude



# Azimuthal Jet Asymmetry At HERA

ZEUS, Physics Letters B 551 (2003) 226-240:

$$\frac{d\sigma}{d\phi_{jet}^{BF}} = A + B \cos(\phi_{jet}^{BF}) + C \cos(2\phi_{jet}^{BF})$$

Azimuthal asymmetry in  
jet-production, at LO and  
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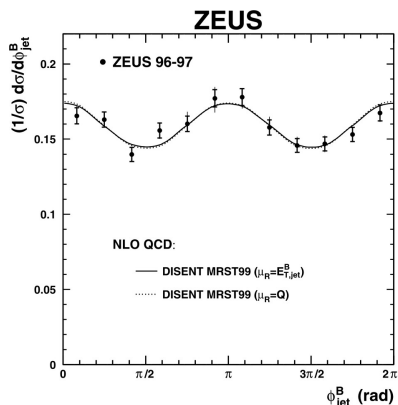
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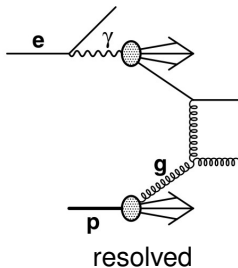
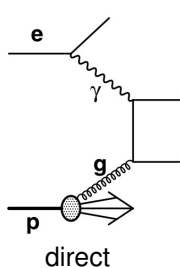
interference between +1 and -1 helicity terms of the transversely polarized part of the exchanged photon

interference between transversal and longitudinal components of the structure function



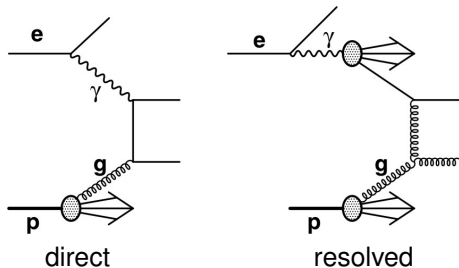
# Dijets In Photoproduction

$$d\sigma_{\gamma p}^{jet} = \sum_{i,j=q\bar{q}g} \int_0^1 \int_0^1 dx_p dx_\gamma f_i(x_p, \mu^2) f_j(x_\gamma, \mu^2) d\hat{\sigma}_{ij}^{jet}(x_p, x_\gamma, \alpha_s, \mu)$$



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define  $\theta^*$  as the dijet scattering angle in the parton-parton centre-of-mass frame,  $\cos \theta^* = \tanh \left( \frac{\eta_1 - \eta_2}{2} \right)$

# Dijets In Photoproduction

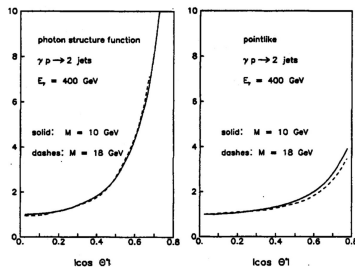
in direct events the dominant boson-gluon fusion processes  $\gamma^* g \rightarrow q\bar{q}$  have a quark propagator

$$\propto \frac{1}{1 - |\cos \theta^*|}$$

in resolved events, the gluon propagator is more prominent,  $qg \rightarrow qg$  and  $gg \rightarrow gg$

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H. Baer et al., Phys. Rev. D 40 (1989) 2844:



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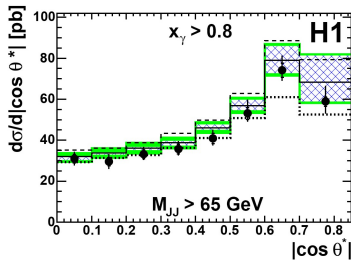
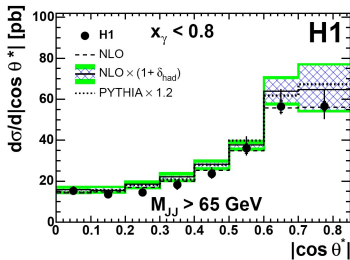
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H1 collaboration, Phys. Lett. B 639 (2006) 21:

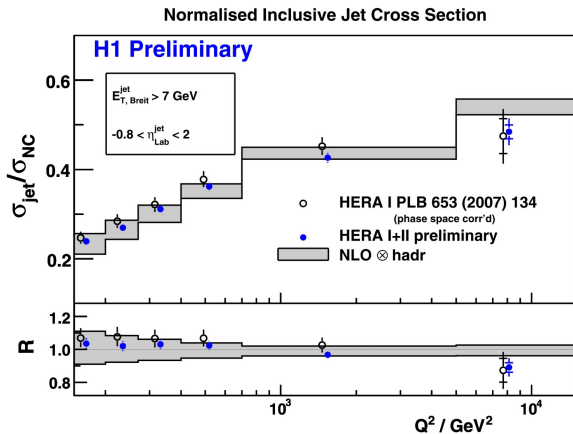


# Precision Physics With Jets



# Determination Of The Strong Coupling

With the full HERA dataset, so-far unachieved precision measurements are possible



tiny experimental errors, partly due to cancellations

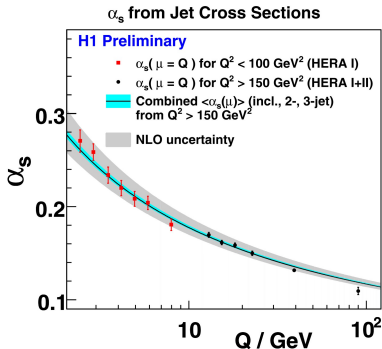
data very well described by fixed-order calculation

theoretical errors dominate in most regions



# Determination Of The Strong Coupling

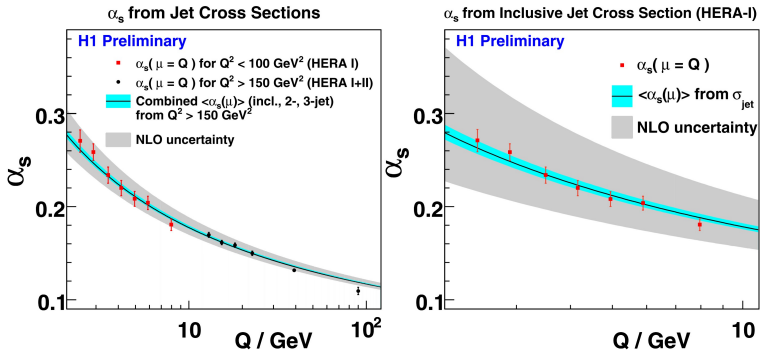
Results from H1prelim-08-31 (black) and H1prelim-08-32 (red)



solid line is  $\alpha_s$  evolved from high  $Q^2$  fits, via two-loop solution of the renormalization group equation

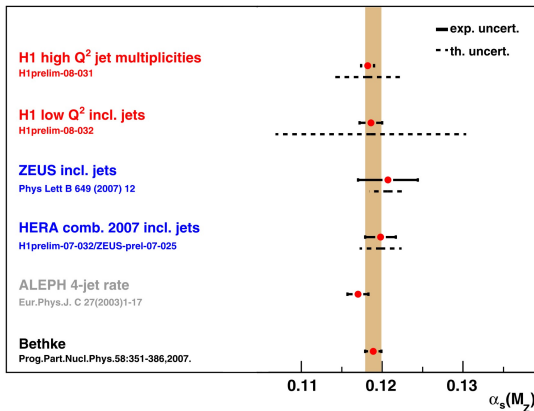
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Results from H1prelim-08-31 (black) and H1prelim-08-32 (red)



solid line is  $\alpha_s$  evolved from high  $Q^2$  fits, via two-loop solution of the renormalization group equation  
 fitting  $\alpha_s$  to low  $Q^2$  data gives huge theoretical errors

# Determination Of The Strong Coupling



experimental errors from HERA jets are competitive with LEP and world average

theoretical errors are dominant because of missing higher orders in NLO calculations

# PDF Determination Including Jets

Parton distribution functions (PDFs) are often determined from inclusive cross sections in global fits within the DGLAP formalism.

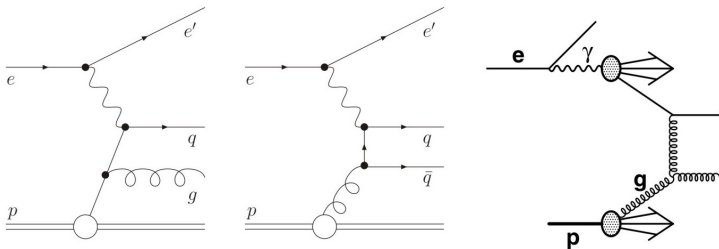
[What about jet data?](#) Reminder:



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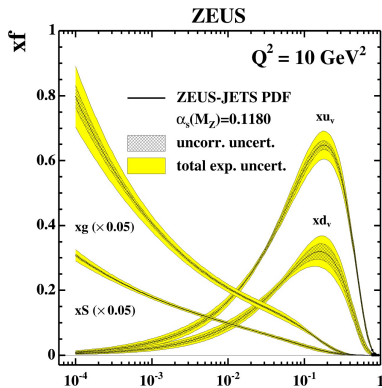
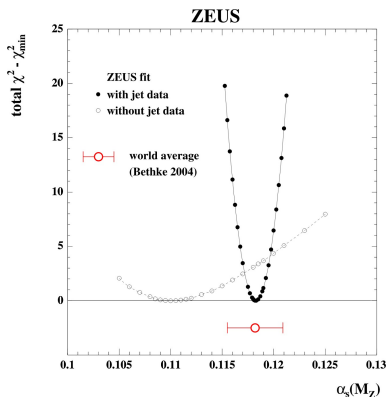
What about jet data? Reminder:



Jet data are directly sensitive to  $\alpha_s$  and the gluon PDF. Instead of using fixed target DIS data for the medium to high  $x$  region, use jet data in the PDF determination.

# PDF Determination Including Jets

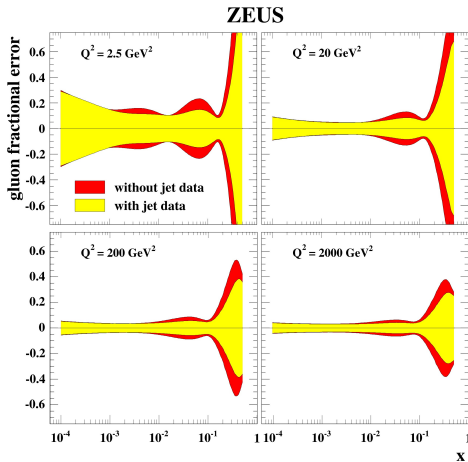
ZEUS collaboration, European Physical Journal C 42 (2005) 1:



Consistent fit without heavy-nuclei corrections, higher-twist contributions or isospin-symmetry assumptions.

# PDF Determination Including Jets

The fitted  $\alpha_s(M_Z)$  is consistent with the world average. The errors on the gluon could be reduced.



The ZEUS-JETS pdf was determined from HERA I data. Dominant errors in global fits are systematic ones, here statistical uncertainties were dominant. Expect improvement for HERA II.

# Summary And Outlook

Jets helped to test and establish QCD.

This talk: jet measurements in the main HERA phase space, where perturbative QCD works very well.

Many topics were not covered,

- + forward jets (small  $x$ , DGLAP?)
- + minijets and multiple interactions
- + jet substructure
- + heavy flavour jets
- + diffraction (factorization?)
- + ...

Wherever pQCD calculations describe data, physical quantities can be extracted. Expect high precision measurements with the full HERA II dataset.

